

Analysis of the Far Ultraviolet Emission Lines in Late Type Stars*

E. de Castro¹, M.J. Fernández-Figueroa¹, M. Rego¹, and D. Ponz²

¹ Departamento de Astrofísica, Facultad de Físicas, Universidad Complutense, Ciudad Universitaria, Madrid-3, Spain

² Villafranca Satellite Tracking Station Apartado 54065, Madrid, Spain

Received December 15, 1980; accepted March 27, 1981

Summary. Far ultraviolet emission lines in a sample of seven stars of spectral type later than F5 and different luminosity class have been analyzed to calculate the electron density and pressure in the formation region of the lines. Three different methods were used. We obtain electron pressures larger than solar for all dwarf stars of the sample. The low resolution observations were made with the IUE satellite.

Key words: emission lines – transition region – corona

1. Introduction

The main source of information about the outer stellar regions of intermediate and late type stars are far ultraviolet observations. The emission lines in this spectral region are associated with chromospheres, transition regions and coronae.

Until quite recently most of our knowledge of outer stellar regions has come from studying the Ca II H and K emission lines in late type stars. Later the Mg II h and k and the He I λ 10830 Å lines were added as indicators of chromospheres, transition regions and coronae.

The IUE satellite now provides a valuable opportunity to study the ultraviolet spectra of cool stars in more detail, and allows to improve the information about the structure of these outer regions.

The presence of emission lines in the spectral range λ 1100–2000 Å indicates a plasma with temperatures larger than 10,000 K. For late type stars one distinguishes between two groups, one with a solar type transition region as well as corona and another with no indication of material hotter than 20,000 K (Linsky and Haisch, 1979). The dwarf stars of this work must be included in the first group.

In this paper we analyze the ultraviolet emission lines in a sample of stars of spectral type later than F5.

2. Observations

A sample of seven stars of different temperatures and luminosity class showing Ca II H and K emission lines has been chosen. The

Send offprint requests to: M.J. Fernández-Figueroa

* Based on observations by the International Ultraviolet Explorer collected at the Villafranca Satellite Tracking Station of the European Space Agency

effective temperature of the atmosphere has been derived from the broad-band color of Johnson et al. (1966) using the mean color- T_{eff} relations of Johnson (1966). The Ca II K intensity (Wilson and Bappu, 1957), the full width at the base (FWB) of Mg II k (Rego et al., 1979; Weiler and Oegerle, 1979) and the equivalent width of He I λ 10,830 Å (Zirin, 1976) are given in Tables 1 and 2.

Table 1. Stellar parameters

HD	Name	Sp. T.	m_v	$B-V$	$V-I$	$V-R$	T_{eff}
3712	α Cas	K0II–III	2.2	1.17	1.38	0.78	4550
4502	ζ And	K1III	4.1	1.10	1.44	0.85	4500
20630	κ Cet	G5V	4.8	0.68	0.93	0.57	5600
62509	β Gem	K0III	1.1	1.00	1.25	0.75	4800
114710	β Com	G0V	4.3	0.58	0.79	0.49	5900
124850	iVir	F7IV	4.1	0.51	0.77	0.50	6200
142373	χ Her	F9V	4.6	0.57	0.80	0.48	6000

Table 2. Chromospheric indicators

HD	I (K CaII)	FWB (k MgII)	W (He)
3712	2	—	0
4502	5	187 km/s ^a	760 mÅ
20630	3	128 km/s ^b	200 mÅ
62509	2	193 km/s ^a	140 mÅ
114710	1	124 km/s ^b	—
124850	1	—	—
142373	1	—	—

^a Weiler and Oegerle (1979) ^b Rego et al. (1979)

Table 3. Observational data

HD	t (min)	Image no.
3712	150	SWP 6697
4502	50	SWP 6834
20630	50	SWP 9462
62509	20	SWP 6472
114710	63	SWP 9465
124850	60	SWP 9464
142373	80	SWP 9463

IUE low resolution ($\approx 6 \text{ \AA}$) observations were obtained in the spectral range $\lambda\lambda 1150\text{--}2000 \text{ \AA}$ with the SWP camera. The exposure time and the image number for each star are given in Table 3. Three stars of the sample αCas , ζAnd , and βGem were not observed by us, but their images were released to the scientific community in June 1980.

Several emission lines of C, O, N, Si in different ionization states have been identified in all spectra. The identified lines, the observed fluxes and the surface fluxes, both in $\text{erg cm}^{-2} \text{ s}^{-1}$, are listed in Tables 4a and 4b for each star. Observed line fluxes have been transformed into surface fluxes with the angular diameter obtained by the relations of Barnes and Evans (1976) as well as Barnes et al. (1976). For αCass it was necessary to correct for

interstellar extinction by means of the ultraviolet interstellar extinction law $\log F_{\lambda}^{\text{cor}} = \log F_{\lambda}^{\text{ob}} + 0.43 k_{\lambda} E(B-V)$, where k_{λ} is the total extinction per unit visual color excess, $E(B-V)$, obtained by Nandy (1976) and Seaton (1979).

The normal short wavelength limit of the IUE SWP spectra is $\lambda 1150 \text{ \AA}$, however this limit was effectively moved up to about $\lambda 1230 \text{ \AA}$ in our observations. The long exposure time required makes the geocoronal $\text{Ly}\alpha$ quite strong, disturbing any feature between $\lambda 1200$ and $\lambda 1230 \text{ \AA}$. On the other hand, the long wavelength limit is approximately $\lambda 1830 \text{ \AA}$ for the hotter stars of the sample due to the absorption edge of Si I.

An upper limit of 3σ has been included in the fluxes listed in Table 4 when the line was not observed.

Table 4a and b. Observed and surface fluxes both in $\text{erg cm}^{-2} \text{ s}^{-1}$

Table 4a

$\lambda(\text{\AA})$	El	HD 20630		HD 114710		HD 124850		HD 142373	
		F_{ob}	F_{*}	F_{ob}	F_{*}	F_{ob}	F_{*}	F_{ob}	F_{*}
1240	N V	6.22 (-14)	9.81 (3)	6.23 (-14)	8.13 (3)	6.19 (-14)	6.06 (3)	5.42 (-14)	9.09 (3)
1304	O I	7.06 (-14)	1.11 (4)	2.34 (-14)	3.05 (3)	1.60 (-13)	1.56 (4)	5.44 (-14)	9.12 (3)
1335	C II	1.37 (-13)	2.17 (4)	3.25 (-14)	4.23 (3)	2.42 (-13)	2.63 (4)	4.21 (-14)	7.06 (3)
1357	O I	2.72 (-13)	4.30 (4)	1.79 (-14)	2.34 (3)	7.95 (-14)	7.77 (3)	4.10 (-14)	6.88 (3)
1394	Si IV	7.66 (-14)	1.21 (4)	2.37 (-14)	3.09 (3)	1.27 (-13)	1.24 (4)	4.45 (-14)	7.47 (3)
1403	Si IV	7.04 (-14)	1.11 (4)	2.96 (-14)	3.86 (3)	9.81 (-14)	9.60 (3)	4.09 (-14)	6.83 (3)
1482	S I	5.96 (-14)	9.41 (3)	3.66 (-14)	4.79 (3)	1.09 (-13)	1.06 (4)	8.33 (-14)	1.39 (4)
1530	Si II	6.65 (-14)	1.99 (4)	7.28 (-14)	9.50 (3)	3.03 (-13)	2.96 (4)	8.99 (-14)	1.51 (4)
1550	C IV	2.17 (-13)	2.47 (4)	7.14 (-14)	9.32 (3)	3.43 (-13)	3.36 (4)	9.39 (-14)	1.58 (4)
1640	He II	6.00 (-14)	9.47 (3)	3.06 (-14)	4.00 (3)	2.76 (-13)	2.70 (4)	—	—
1657	C I	1.47 (-13)	2.31 (4)	8.07 (-14)	1.05 (4)	1.90 (-13)	2.84 (4)	9.94 (-14)	1.67 (4)
1748	Ni II	1.33 (-13)	2.10 (4)	—	—	—	—	—	—
1808	Si II	2.19 (-13)	3.45 (4)	1.65 (-13)	2.15 (4)	sat	sat	<3.36 (-13)	<5.63 (4)
1817	Si II	2.87 (-13)	4.53 (4)	1.92 (-13)	2.51 (4)	sat	sat	<4.35 (-13)	<7.31 (4)
1892	Si III	2.52 (-13)	3.97 (4)	<3.37 (-13)	<4.04 (4)	sat	sat	sat	sat
1909	C III	3.02 (-13)	4.77 (4)	<3.67 (-13)	<4.35 (4)	sat	sat	sat	sat

Table 4b

$\lambda(\text{\AA})$	El	HD 3712		HD 4502		HD 62509	
		F_{ob}	F_{*}	F_{ob}	F_{*}	F_{ob}	F_{*}
1240	N V	2.08 (-13)	9.04 (2)	2.49 (-13)	9.31 (3)	1.78 (-13)	4.50 (2)
1304	O I	1.09 (-12)	4.73 (3)	9.74 (-13)	3.62 (4)	5.52 (-13)	1.46 (3)
1335	C II	—	—	3.74 (-13)	1.16 (4)	1.80 (-13)	1.46 (3)
1357	O I	—	—	9.55 (-14)	3.58 (3)	—	—
1394	Si IV	1.14 (-14)	4.99 (2)	2.50 (-13)	8.40 (3)	1.08 (-13)	3.19 (2)
1403	Si IV	2.43 (-13)	1.05 (3)	1.98 (-13)	7.38 (3)	1.04 (-13)	2.76 (2)
1530	Si II	3.20 (-13)	1.34 (3)	1.11 (-13)	4.14 (3)	—	—
1550	C IV	<2.29 (-13)	<1.00 (3)	5.51 (-13)	2.06 (4)	2.03 (-13)	5.34 (2)
1640	He II	—	—	2.38 (-13)	8.88 (3)	1.74 (-13)	4.42 (2)
1657	C I	2.26 (-13)	9.86 (2)	2.37 (-13)	8.86 (3)	2.95 (-13)	7.71 (2)
1748	Ni II	—	—	1.13 (-13)	4.22 (3)	—	—
1808	Si II	<2.50 (-13)	<1.32 (3)	2.79 (-13)	1.04 (4)	3.89 (-13)	1.03 (3)
1817	Si II	<4.72 (-13)	<2.05 (3)	6.12 (-13)	2.28 (4)	5.27 (-13)	1.38 (3)
1892	Si III	<3.13 (-13)	<1.36 (3)	2.67 (-13)	9.96 (3)	<4.04 (-13)	<1.06 (3)
1909	C III	<3.30 (-13)	<1.43 (3)	1.90 (-13)	7.11 (3)	<3.60 (-13)	<9.41 (2)

3. Results and Discussion

The ultraviolet spectra of all dwarf stars, namely κ Cet, β Com, and χ Her, and of the subgiant iVir appear clearly similar to that of the quiet sun with typical lines of the solar transition region and corona. The more luminous stars β Gem, ζ And, and α Cas show different spectra than the solar one, however it would not be correct to include these three stars into a single group, as for instance, high-temperature ions like C IV are present in β Gem but not in α Cas. On the other hand, the ultraviolet spectrum of the most luminous star, ζ And, has similar lines and comparable surface fluxes to that found in solar active regions, this result is not surprising since ζ And is a RS CVn star.

Figure 1 shows the H-R diagram with boundaries for circumstellar lines (Reimers, 1977), Mullan's (1978) supersonic transition locus, and the sharp empirical line of Linsky and Haisch (1979). We have located our stars on this diagram. Dwarf stars, the giant β Gem and the subgiant iVir are shown to the left of all boundary lines, they are "solar type stars", while ζ And is immediately to the right of the Linsky and Haisch dividing line, although this line may not be as clearly defined as it appears in the figure (Linsky and Haisch, 1979). With respect to α Cas, lies above all the boundaries and its spectrum shows rather weak emission lines and is very similar to spectra of "non solar type" stars. For this reason, α Cas may be included in the supersonic wind class while the rest of the stars in our sample can be included in the transonic wind or coronal wind class.

The ratio of the C IV, Si IV and N V surface fluxes to the total surface flux are plotted versus $(B - V)$ in Fig. 2. All dwarf stars are more active than the sun and, apparently, there is no dependence on the effective temperature.

The number of more luminous stars is not enough to infer a dependence on the surface gravity, however the selected sample shows a similar scatter than that found by Linsky et al. (1978) from the Ca II line fluxes.

By assuming that the emission lines in the transition region and corona are optically thin and formed by collisional excitation from

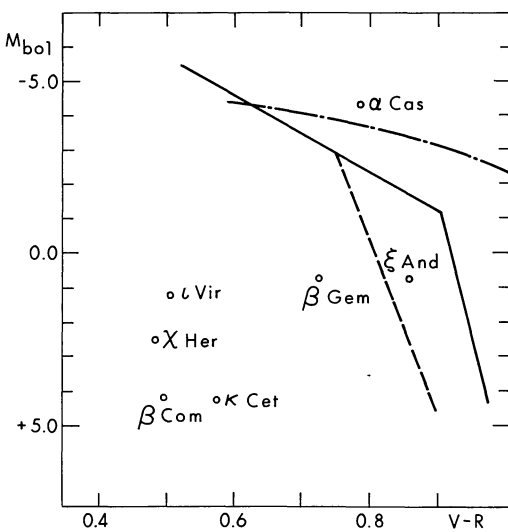


Fig. 1. H-R diagram for stars of our sample. The solid line is Mullan's, the dashed is from Haisch and Linsky and the dot-dashed is from Reimers

the ground state, the peak formation temperatures T_{\max} for each line have been calculated, where T_{\max} is the temperature at which the function $g(T)$ is a maximum,

$$g(T) = T_e^{-1/2} 10^{-5040W/T} \frac{N(\text{ion})}{N(\text{el})}$$

T_e refers to electron temperature, W is the excitation potential, and $N(\text{ion})/N(\text{el})$ tabulated by Jordan (1969). According to these values the emission lines can be divided into three groups,

- $T < 2 \cdot 10^4$ K C I, O I lines
- $2 \cdot 10^4 < T < 7 \cdot 10^4$ K C II, Si II lines
- $T > 7 \cdot 10^4$ K N V, C IV, Si IV lines

In all stars of our sample there are lines corresponding to each of these groups. The lines within group *c* are indicators of transition regions and possibly coronae, unfortunately the geocoronal Ly α line masks the identification of O V and O VI lines the formation temperature of which is about 10^6 K and could be proof of the existence of a corona.

The electron density in the region where ultraviolet emission lines are formed, has been calculated using three different methods.

1. Relative Intensities of Intercombination Lines

According to Doschek et al. (1978), the slope of the emission measure, $A_{\text{el}} N_e^2 V$, plotted as a function of temperature is the same in all stars. The line ratios C IV/Si IV and N V/Si IV are given in Fig. 3, for all stars in our sample and for stars included in the works of Linsky et al. (1978) and Jordan (1980). It is clear from the Fig. 3 that the conclusion of Doschek et al. (1978) is only valid for stars with the same luminosity class, this can be explained by the fact that the line ratio depends on density through the quantity $N(\text{ion})/N(\text{el})$ which is not the same for stars with different luminosity class.

If we assume the approximation proposed by Doschek et al. (1978), for electron densities lower than 10^{11} cm^{-3} , the line ratio C III ($\lambda 1909$)/Si III ($\lambda 1892$) is proportional to N_e^{-1} . Using the

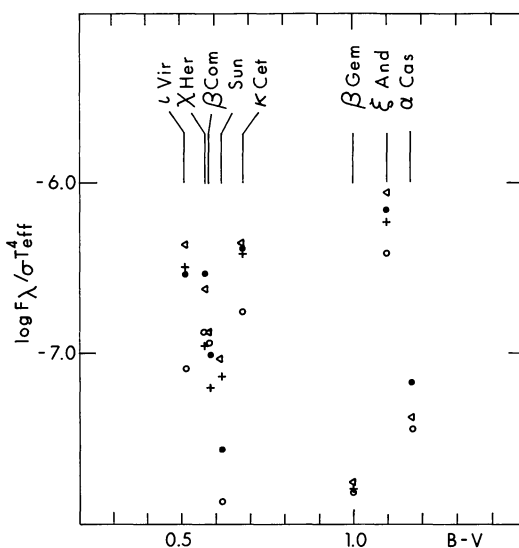


Fig. 2. Ratio surface flux to the total surface luminosity versus $(B - V)$, for C II $\lambda 1335 \text{ \AA}$ line (+), C IV $\lambda 1550 \text{ \AA}$ (Δ), Si IV $\lambda 1403 \text{ \AA}$ (\circ) and N V $\lambda 1240 \text{ \AA}$ (\bullet)

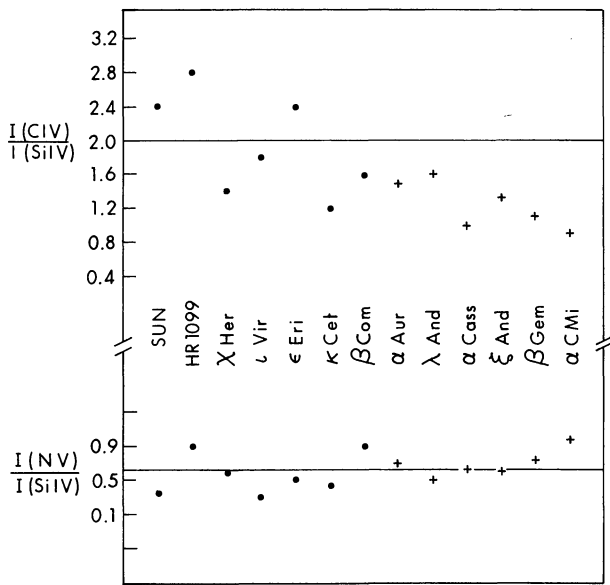


Fig. 3. Flux ratios of C IV (1550) and N V (1240) relative to Si IV (1403 + 1394) for the stars indicated. Filled points: luminosity class V, crosses: other luminosity classes

theoretical values deduced for the sun and the observed ratio for κ Cet we find an electron pressure $P_e = 8.32 \cdot 10^{14} \text{ cm}^{-3} \text{ K}$.

A similar method has been proposed by Cook and Nicolas (1979) using the lines $\lambda 1402$ Si IV and $\lambda 1909$ C III from which we obtain $P_e = 9 \cdot 10^{14} \text{ cm}^{-3} \text{ K}$.

In our sample the only star with C III and Si III lines and solar luminosity class is κ Cet, for this reason the method is only applied to this star.

2. Energy Balance in the Transition Region

Assuming the steady-state local energy equation in which the divergence of the downward conductive flux is balanced by radiative losses and the collisional excitations are exactly balanced by radiative de-excitations, the surface flux ratios in the vicinity of T_{max} for each ion are given by,

$$\frac{F(*)}{F(\odot)} = \frac{A(\text{el})_* P_*}{A(\text{el})_{\odot} P_{\odot}}$$

where A is the element abundance (Haisch and Linsky, 1976). The above relationship has been used to calculate the pressures listed in Table 5 for lines with peak temperature larger than $2 \times 10^4 \text{ K}$. In all dwarf stars, the calculated pressure is larger than solar. Assuming $P_{\odot} = 0.15 \text{ dyn cm}^{-2}$ we obtain $P_e = 10^{16} \text{ cm}^{-3} \text{ K}$ for κ Cet, that is 10^2 times larger than the value obtained previously.

3. Cool Corona

Several authors have suggested the existence of coronae with temperature lower than 10^6 K . Coronal line fluxes can be determined assuming that the line formation takes place in a plane - parallel isothermal region and are given by the relationship (McClintock et al., 1975),

$$F(\lambda) = \frac{7.07 \cdot 10^{25} \Omega}{g_1 \lambda T_c^{3/2}} A(\text{el}) \frac{N(\text{ion})}{N(\text{el})} \exp(-W/kT_c) \frac{P^2}{g}$$

Table 5. Pressures obtained from the "energy balance" method

		P_*/P_{\odot}			
		1240	1335	1394	1550
HD	3712	1.5	—	0.9	0.2
HD	4502	17.3	4.0	7.8	5.6
HD	20630	9.6	4.0	6.2	5.0
HD	62509	0.3	0.1	0.1	0.1
HD	114710	5.0	0.5	4.8	0.6
HD	124850	19.0	1.8	18.0	4.0
HD	142373	19.0	4.0	17.0	6.0

Table 6. Total pressures obtained from the "cool corona" method

		$P_{\text{min(cool corona)}} \text{ dyn/cm}^2$
HD	3712	1.47 (-3)
HD	4502	2.2 (-3)
HD	20630	4.7 (-2)
HD	62509	1.0 (-3)
HD	114710	3.3 (-2)
HD	124850	2.5 (-2)
HD	142373	3.8 (-2)

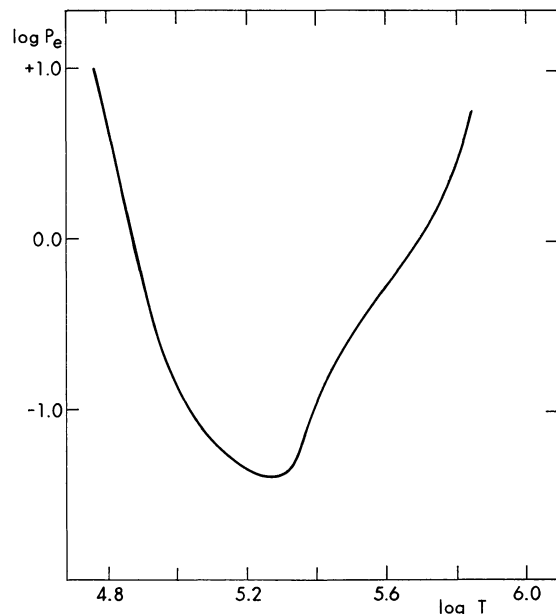


Fig. 4. Pressure as a function of temperature for N V 1240 Å line of κ Cet

where Ω is the collision cross-section, g_1 the statistical weight and g the surface gravity. In Fig. 4 the pressure, P , of κ Cet has been plotted as a function of temperature for N V 1240 Å. This curve shows the possible values of pressure for a corona of this type, the

minimum value corresponding to T_{\max} . The method provides a lower limit for the transition region pressure. Table 6 gives this lower limit for all the stars in our sample.

Errors in the computation of electron pressures using the above three methods are due to the uncertainty of the observed fluxes and are less than 20%. Errors due to uncertainties in atomic data or to the assumed hypotheses are not considered. The validity of our hypotheses are discussed in the conclusions.

4. Conclusions

All the stars in our sample show transition region and corona indicators similar to the sun except for α Cas. The analyzed dwarf stars are more active than the sun.

For all observed dwarf stars the electron pressures obtained by each of three methods used are greater than the corresponding solar values.

The pressure of the giant β Gem is lower than the dwarf stars. This suggests a trend of decreasing pressure with increasing stellar luminosity.

The line surface fluxes in ζ And are higher by a factor of ≥ 10 than in the sun and the calculated pressures are similar to a solar active region. This result is not opposed to the above suggestion since ζ And is a RS CVn star.

Our calculations are only approximate and the following limitations of the methods used should be noted,

The "relative intensities of intercombinations lines" method is only valid for solar luminosity class stars.

The "energy balance" method only defines an upper limit for the pressure and we have neglected the acoustic and magnetic fluxes which should perhaps be included.

From the "cool corona" method we have obtained a lower limit to the pressure in the transition region but, again, neglecting the acoustic and magnetic pressures.

These limits on electron pressure will be used as starting point in the development of a more elaborate model of the transition region.

References

- Barnes, J.G., Evans, D.S.: 1976, *Monthly Notices Roy. Astron. Soc.* **174**, 489
- Barnes, J.G., Evans, D.S., Parsons, S.B.: 1976, *Monthly Notices Roy. Astron. Soc.* **174**, 503
- Brown, A., Jordan, C.: 1980, Proceedings of the Second European IUE Conference. Tübingen, Germany
- Cook, J.W., Nicolas, K.R.: 1979, *Astrophys. J.* **229**, 1163
- Doschek, G.A., Feldman, U., Mariska, J.T.: 1978, *Astrophys. J. Letters* **226**, L35
- Haisch, B.M., Linsky, J.L.: 1976, *Astrophys. J. Letters* **205**, L39
- Johnson, H.L.: 1966, *Ann. Rev. Astron. Astrophys.* **4**, 193
- Johnson, H.L., Mitchell, R.I., Iriarte, B., Wisniewski, W.Z.: 1966, *Comm. Lunar Planetary Lab.* **4**, 99
- Jordan, C.: 1969, *Monthly Notices Roy. Astron. Soc.* **142**, 501
- Kelch, W.J., Linsky, J.L., Wordn, S.P.: 1979, *Astrophys. J.* **229**, 700
- Linsky, J.L., Haisch, B.M.: 1979, *Astrophys. J. Letters* **229**, L27
- Linsky, J.L., Wordn, S.P., McClintock, W., Robertson, R.M.: 1979, *Astrophys. J. Suppl.* **41**, 47
- Linsky, J.L., Ayres, T.R., Basri, G.S., Morrison, N.D., Boggess, A., Schiffer III, F.H., Holm, A., Cassatella, A., Heck, A., Macchetto, F., Stickland, D., Wilson, R., Blanco, C., Dupree, A.K., Jordan, C., Wing, R.F.: 1978, *Nature* **275**, 389
- McClintock, W., Linsky, J.L., Henry, R.C., Moos, H.W., Gerola, H.: 1975, *Astrophys. J.* **202**, 165
- Mullan, D.J.: 1978, *Astrophys. J.* **226**, 51
- Nandy, K., Thompson, G.I., Jamar, C., Monfils, A., Wilson, R.: 1976, *Astron. Astrophys.* **51**, 63
- Pottasch, S.R.: 1964, *Space Sci. Rev.* **3**, 816
- Rego, M., Fernández-Figueroa, M.J.: 1979, *Astron. Astrophys.* **76**, 249
- Reimers, D.: 1977, *Astron. Astrophys.* **57**, 395
- Seaton, M.J.: 1979, *Monthly Notices Roy. Astron. Soc.* **187**, 73
- Seaton, M.J.: 1962, Atomic and molecular processes. Academic Press. New York
- Weiler, E.J., Oegerle, W.R.: 1979, *Astrophys. J. Suppl.* **39**, 537
- Wilson, O.C., Bappu, M.K.V.: 1957, *Astrophys. J.* **125**, 661
- Zirin, H.: 1976, *Astrophys. J.* **208**, 414